Satellite Test of the Isotrophy of the One-Way Speed of Light Using ExTRAS

Peter Wolf

Bureau International des Poids et Mesures,
Pavillon de Breteuil, 92312 Sèvres CEDEX, France
and
Queen Mary and Westfield College, School of Mathematical Sciences,
Mile End Road, London E1 4NS, Great Britain

Abstract: A test of the second postulate of special relativity, the universality of the speed of light, using the ExTRAS (Experiment on Timing Ranging and Atmospheric Sounding) payload to be flown on board a Russian Meteor-3M satellite (launch date January 1997) is proposed. The propagation time of a light signal transmitted from one point to another without reflection would be measured directly by comparing the phases of two hydrogen maser clocks, one on board and one on the ground, using laser or microwave time transfer systems. An estimated uncertainty budget of the proposed measurements is given, resulting in an expected sensitivity of the experiment of $\delta c/c < \delta x 10^{-10}$ which would be an improvement by a factor of ~ 430 over previous direct measurements and by a factor of ~ 4 over the best indirect measurement. The proposed test would require no equipment additional to what is already planned and so is of inherently low-cost. It could be carried out by anyone having access to a laser or microwave ground station and a hydrogen maser.

1. Introduction

Einstein's second postulate, affirming the universality of the speed of light for inertial frames, is fundamental to the theories of special and general relativity. It can be tested directly by comparing the propagation times of two light signals travelling from one point to another along the same path but in opposing directions (often referred to as a test of the isotropy of the one-way speed of light). The only such test, was carried out by Krisher

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et al. [1], who compared the phases of two hydrogen masers separated by a distance of 21 km and linked via an ultrastable fibre optics link of the NASA deep space network. The sensitivity of this experiment, expressed as a limit on the anisotropy of the speed of light, was $\delta c/c < 3.5 \times 10^{-7}$, where c is the velocity of light in vacuum. Riis et al. [2] tested the isotropy of the first order Doppler shift of light emitted by an atomic beam (and indirectly thereby the second postulate) using fast-beam laser spectroscopy obtaining the currently best limit on the anisotropy, $\delta c/c < 3 \times 10^{-9}$. This presents a 10 fold improvement on previous values from experiments measuring the isotropy of the first order Doppler shift using the frequency links in the NASA GP-A rocket experiment [3] and so-called Mössbauer rotors [4, 5]. In the test theory of Mansouri and Sexl [6] the above results can be interpreted as limits on the parameter α using the relation $\delta c/c = (1+2\alpha)v/c$ [1, 2] where v is the velocity of the Earth with respect to the mean rest frame of the universe ($v \sim 300$ km/s). This yields values of $\alpha = -1/2 \pm 1.8 \times 10^{-4}$ and $-1/2 \pm 1.4 \times 10^{-6}$ for the experiments by Krisher et al. [1] and Riis et al. [2] respectively.

The experiment proposed here would test the isotropy of the transmission time of light signals between two points directly and on a non-laboratory scale with an estimated accuracy of $\delta c/c < 8 \times 10^{-10}$, using the ExTRAS payload on board the Russian Meteor-3M satellite scheduled for launch in January 1997, T2L2 (Time Transfer by Laser Light) time transfer and a hydrogen maser at the ground station. This, if realized, would present a 430 fold improvement on previous direct measurements [1] and a slight improvement on the value obtained by Riis et al. [2]. In Section 2 the principle of the experiment is explained while Section 3 provides an evaluation of its sensitivity aimed at including all error sources that may exceed one picosecond and based on the uncertainty budget for the T2L2 method by Thomas & Uhrich [7].

2. Experimental principle

The ExTRAS payload consists of two active, auto-tuned hydrogen masers communicating with ground stations via a PRARE (Precise Range and Range-Rate Equipment) microwave link and a T2L2 laser link. Once operational, the system should reflect laser pulses, emit and receive microwave signals and date all such events on the on-board time scale provided by the hydrogen masers. The satellite will follow polar orbit, at an altitude of 1000 km with a period of order 100 min and a duration of one passage of ~ 17 min.

In principle, the proposed experiment is similar to that performed by Krisher et al. [1]. A laser signal emitted from the station E is reflected at the satellite S and returned to E (see figure 1). The readings of the ground hydrogen maser at emission (τ_0) and reception (τ_2) and that of the space maser at the moment of reflection (τ_1) are recorded. The differences $\tau_1^{-\tau_0}$ and $\tau_2^{-\tau_1}$ represent the up and down transmission times T_1 and T_2 respectively plus some initial phase difference of the clocks. Note that no synchronization convention or

procedure is assumed. Einstein's second postulate would require that for a series of measurements, after accounting for the path asymmetries, the difference T_1 - T_2 should be equal to a constant Δ_0 (due to the initial clock offset) independent of the spatial orientation of the individual links. More particularly one obtains for a single link (see [8] for more detail),

$$T_1 - T_2 = \Delta_0 + 2 \vec{R}(t_e) \cdot \vec{v}(t_e) / c^2 + (\Delta_{i(up)} - \Delta_{i(down)}) + O(c^{-3})$$
 (1)

where $\overrightarrow{R}(t_e)$ is the vector from E to S at the coordinate time of emission of the signal t_e in a geocentric, inertial reference frame, $\overrightarrow{v}(t_e)$ is the velocity of the ground station at signal emission in the same frame and Δ_i are internal delays (cables etc.).

The initial clock offset Δ_0 is a constant, provided that the two clocks are syntonized. This can be achieved at the 10^{-15} accuracy level (the best hydrogen maser stability) using time transfer data over a sufficiently long integration period and taking into account all known effects (gravitational redshift, second order Doppler, maser drift). One would expect the effect on the syntonization, of an eventual anisotropy of the propagation time of the light signals, to average out in a global treatment using time transfers in all spatial directions.

Terms of order c^2 amount to ~ 40 ns and can be calculated to picosecond accuracy if $R(t_e)$ and $v(t_e)$ are known to within ~ 50 m and ~ 0.01 m/s respectively, which represents no difficulty for modern satellite orbitography. Of course, a possible anisotropy would also have an effect on the satellite orbit determination, but as the range R cancels to first order in (1) this effect would be negligible. Furthermore, the satellite orbit is obtained from round-trip ranging measurements, which should, again to first order, be insensitive to anisotropy of the propagation time of the light signals.

Terms of order c^{-3} can amount to several picoseconds but can be calculated to picosecond accuracy without difficulty [8]. The effect of asymmetry in the atmospheric delays for the up and down links is below one picosecond.

Hence, after accounting for path asymmetry, any variation of the difference $T_{\rm l}$ - $T_{\rm l}$ with the spatial orientation of the laser link should be due to a violation of the second postulate.

3. Estimation of the experiment sensitivity

The sensitivity of the proposed test can be estimated by considering two individual laser links as shown in figure 2. The time intervals $\tau_2^-\tau_0^-$, $\tau_5^-\tau_3^-$ and $\tau_3^-\tau_0^-$ are measured using the ground hydrogen maser with the interval $\tau_4^-\tau_1^-$ obtained from the space hydrogen maser. Designating the individual transmission times by T_1^- , T_2^- , T_3^- and T_4^- as shown in figure 2 and

assuming that one of the links is colinear with the direction of the presumed anisotropy, the difference between the two links is given by,

$$(T_1 - T_2) - (T_3 - T_4) + \Delta_s = 2 \Delta_s (1 - \cos\theta).$$
 (2)

Here Δ_s represents the correction due to the path asymmetries of the individual links; Δ_a is the maximum delay for a single transmission due to the anisotropy, and θ is the angle between the two links in the inertial geocentric frame.

If Einstein's second postulate is true the right hand side of equation (2) should be equal to zero within the measurement error.

The experiment should be capable of detecting an anisotropy under the condition

$$\varepsilon < 2 \Delta_{R}(1-\cos\theta),$$
 (3)

where ε represents the total measurement uncertainty.

The sensitivity of the experiment is therefore given by,

$$\delta c/c = \Delta_{\bullet}/T = \varepsilon/[2T(1-\cos\theta)] \tag{4}$$

where T is a typical transmission time $(T_{max} \sim 12 \text{ ms})$.

Maximal sensitivity is achieved when the measurements are taken at the beginning and the end of a single passage of the satellite directly above the station. In this case $\theta \sim 180^{\circ}$, $T \sim 12$ ms and the error accumulated due to the instability of the hydrogen masers is very small because of the short integration time of ~ 17 min. Table 1 lists the individual sources of uncertainty that are estimated to exceed 1 ps. Four sources of uncertainty are listed in the table:

- (i) The stability of the hydrogen masers for integration times of 1000 s is of the order 2.1 parts in 10^{15} [7] which gives an accumulated uncertainty of ~ 2 ps per maser over an integration time of 17 min.
- (ii) As systematic errors in the on-board payload cancel when the two links are differenced, only its instability over 17 min contributes. Ten picoseconds [7] seems a conservative estimate for such a short integration time.
- (iii) Only the instability of the Earth station during the experiment contributes. Degnan [9] states that the precision of satellite laser ranging stations is of order 1 to 3 mm, which corresponds to an uncertainty of less than ten picoseconds.

(iv) Information on the counter uncertainties is provided by the T2L2 proposing team.

In the calculation of $(T_1 - T_2) - (T_3 - T_4)$ the differences $\tau_4 - \tau_1$ and $\tau_3 - \tau_0$ measured by the space and ground clock respectively appear with a factor of 2. Hence all uncertainty sources participating in the measurement of these intervals ((i),(ii),(iv)) have been multiplied by this factor.

For the measurement of anisotropy in a direction which is not in the plane of orbit, the two links are separated by the time necessary for the Earth station to change its position with the rotation of the Earth so as to see the satellite from opposing directions (~ 14500 s). The hydrogen maser stability for such integration times is of the order 1,5 parts in 10^{15} [7], which gives an uncertainty of $\sim 44\sqrt{2}$ ps in (i). Contributions from other error sources are those given in Table 1. Hence the value for the total measurement uncertainty is $\epsilon \sim 72$ ps. Note also that in this case θ cannot exceed 120° .

Substituting these values for ε and θ into (4) gives an experimental sensitivity of $\delta c/c = 7.9 \times 10^{-10}$ when the direction of the anisotropy lies in the orbital plane of the satellite and $\delta c/c = 2 \times 10^{-9}$ otherwise. Following Krisher et al. [1] the experiment can be interpreted in the framework of the test theory by Mansouri & Sexl [6] resulting in limits on the parameter α of $\alpha = -1/2 \pm 4 \times 10^{-7}$ and $\alpha = -1/2 \pm 1 \times 10^{-6}$ for the two cases, assuming v = 300 km/s.

Conclusion

The proposed test of the special theory of relativity is expected to improve the upper limit on anisotropy of the propagation time of light signals obtained from the best previous direct measurement [1] by a factor of ~ 430 . It should also provide an improvement (by a factor of ~ 4) on the value inferred from the measurement of the first order Doppler shift by Riis et al. [2]. The extension of this type of experiment to space-time domains (separation of the clocks of ~ 3700 km) which are not attainable in a laboratory may also be an advantage. And last but not least, the experiment does not call for the insallation of additional equipment, hence it can be considered an essentially no-cost experiment which is generally a decisive factor for research in fundamental science.

The same experiment could be performed using the PRARE microwave transfer system in the two-way ranging mode [7] rather than the T2L2 links. This might be of advantage as the PRARE method is not weather dependent. However, uncertainties in the ionospheric propagation delays due to different up and down link frequencies introduce an additional uncertainty of ~ 20 ps per link, which slightly decreases the overall sensitivity of the

experiment to $\delta c/c = 9.8 \times 10^{-10}$ for the case where the direction of the anisotropy lies in the orbital plane of the satellite and to $\delta c/c = 2.1 \times 10^{-9}$ otherwise.

It is likely that the sensitivity of the experiment can be improved if data taken continuously during the passage of the satellite is used to search for the sinusoidal variation with θ of the signal due to anisotropy. Furthermore, if a likely orientation of the presumed anisotropy is identified, for example the direction of the observed dipole anisotropy of the cosmic microwave background [10], it should be possible to improve the experimental sensitivity by statistical treatment of data from different stations and from repeated measurements.

Finally, it should be mentioned that the same type of experiment would yield increased accuracy if performed on satellites at higher altitudes, as this would decrease the ϵ/T ratio in (4). One possible candidate is the Radioastron 1 mission (apogee 85000 km, perigee 2000 km) scheduled for launch in late 1996.

Acknowledgements: Helpful discussions with Dr. Claudine Thomas and Gérard Petit are gratefully acknowledged.

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Tables

Source of uncertainty	σ/ps	
Hydrogen masers(i)	4 √2	
On-board payload(ii)	20	
Earth station(iii)	10 √2	
Counters(iv)	20 √2	
Total(quadratic sum)	ε = 38	

Table 1: Anticipated uncertainty budget for measurement of an anisotropy whose direction lies in the orbital plane. All uncertainties are in picoseconds and correspond to an estimated one standard uncertainty, s.

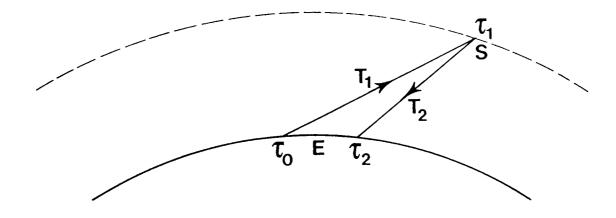


Figure 1: Two-way laser link between an Earth station and the satellite viewed in a geocentric, inertial frame.

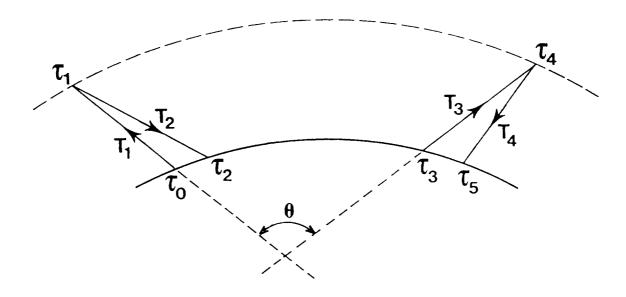


Figure 2: A pair of two-way laser links between an Earth station and the satellite, as viewed in a geocentric, inertial frame.

QUESTIONS AND ANSWERS

LUTE MALEKI (JPL): The experiment that we did at JPL, as you know, was limited because of differential drift of the two H-masers which are not deterministic.

PETER WOLF (BIPM): Yeah, that was the first line of the error budget, which was — that is what I meant by "hydrogen maser," their instability over the integration time, just to accumulate an error in time.

LUTE MALEKI (JPL): No, I'm not talking about the individual instability, I'm talking about the drift that is indeterministic; one maser moves on way, and the other maser moves the other way.

PETER WOLF (BIPM): I didn't consider that. I will have to look into that. Thank you anyway.

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OFFICIAL ATTENDEES' LIST

Khalid S. Al-Dawood

Saudi Arabian Standards Organization P.O. Box 3437 11471 Riyadh SAUDI ARABIA

Khalid K. Al-Dossary

Saudi Arabian Standards Organization P.O. Box 3437 11471 Riyadh SAUDI ARABIA

David W. Allan

Allan's Time P.O. Box 66 Fountain Green, UT 84632 USA

John C. Arnold

AlliedSignal Technical Services Corporation One Bendix Road Columbia, MD 21245 USA 410/964-7931

Franklin G. Ascarrunz

University of Colorado 248 Pheasant Run Louisville, CO 80027 USA

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TASC

1190 Winterson Road Linthicum, MD 21090 USA 410/850-0070

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Francoise S. Baumont

Observatoire de la Cote d'Azur CERGA Avenue Nicolas Copernic 06130 Grasse FRANCE 33 93 405338

Henry V. Bazak, Jr.

Hughes Space and Communication P.O. Box 92919 Los Angeles, CA 90274 USA 310/416-4179

Ronald L. Beard

U.S. Naval Research Laboratory Code 8150 4555 Overlook Avenue, Southwest Washington, DC 20375-5000 USA 202/767-2595

G. Thomas Becker

Air System Technologies, Incorporated 14232 Marsh Lane, Suite 339 Dallas, TX 75234 USA 214/402-9660

Jacques Beser

3S Navigation 23141 Plaza Ptc Drive Laguna Hills, CA 92653 USA

Robert E. Blair

U.S. Air Force Newark Air Force Base Newark, OH 43057 USA 614-522-7377

Scott S. Blake

NASA/GSFC Code 5313 Greenbelt, MD 20710 USA 301/286-3224

Richard E. Blumberg

U.S. Naval Observatory 3450 Massachusetts Avenue, Northwest Washington, DC 20392-5420 USA 202/653-1538

Jimmie Brad

U.S. Naval Research Laboratory Code 8150 4555 Overlook Avenue, Southwest Washington, DC 20375-5000 USA 202/767-5039

Keith L. Brandt

Spectrum Geophysical Instruments 1900 West Garvey Avenue South Suite 200 West Covina, CA 91790 USA 714/544-3000

Lee A. Breakiron

U.S. Naval Observatory 3450 Massachusetts Avenue, Northwest Washington, DC 20392-5420 USA 202/653-1888

Gerhard Brenninger

ESG Elektroniksystem-und Logistik GmbH Einsteinstrasse 174 D-81675 Munchen GERMANY 49 89 9216 2681

Ellis H. Bryant, Jr.

Weatherchron Company 5400 New Peachtree Road Atlanta, Georgia 30341 USA

James A. Buisson

Antoine Enterprises 7714 Martel Place Springfield, VA 22152 USA

Edward E. Burkhardt

Buckhardt Monitoring Service P.O. Box 1411 Glen Allen, VA 23060 USA 804/261-1800

James A. Burton

CNO N63 12739 Purdham Drive Woodbridge, VA 22192 USA 703/695-8202

Bob Bush

Leitch, Incorporated 920 Corporate Lane Chesapeake, VA 23320 USA

Edgar W. Butterline

AT&T 900 Routes 202 and 206 Bedminster, NJ 07921 USA 908/234-4545

Malcolm Calhoun

Jet Propulsion Laboratory 4800 Oak Grove Drive Pasadena, CA 91109 USA 818/354-9763

James C. Camparo

The Aerospace Corporation P.O. Box 92957 Los Angeles, CA 90009 USA 310/336-6944

Thomas P. Celano

TASC 12100 Sunset Hills Road Reston, VA 22090 USA 703/834-5000

Harold Chadsey

U. S. Naval Observatory 3450 Massachusetts Avenue, Northwest Washington, DC 20392-5420 USA 202/653-1888

Laura G. Charron

U.S. Naval Observatory 3450 Massachusetts Avenue, Northwest Washington, DC 20392-5420 USA 202/653-1529

Raymond E. Claflin, III

Claflin Associates 122 Morningside Drive Leominster, MA 01453 USA 508/534-4777

Randolph T. Clarke

U.S. Naval Observatory 3450 Massachusetts Avenue, Northwest Washington, DC 20392-5420 USA 202/653-1034

Dwin Craig

U.S. Naval Research Laboratory Code 8150 4555 Overlook Avenue, Southwest Washington, DC 20375-5000 USA 202/404-7060

Leonard S. Cutler

Hewlett-Packard Company 3500 Deer Creek Road Palo Alto, CA 94303 USA 415/857-5259

William B. Dabney

U.S. Naval Observatory 3450 Massachusetts Avenue, Northwest Washington, DC 20392-5420 USA 202/653-1549

Gerrit de Jong

NMi Van Swinden Laboratorium P. O. Box 654 2600AR Delft THE NETHERLANDS 31 15 691 623

Steve D. Deines

Rockwell International 350 Collins Road, Northeast MS 152-264 Cedar Rapids, IA 52498 USA 319/395-5038

Edoardo V. Detoma

FIAT CIEI - SEPA Corso Giulio Cesare 300 10154 Torino ITALY 39 11 268 2523

James A. DeYoung

U.S. Naval Observatory 3450 Massachusetts Avenue, Northwest Washington, DC 20392-5420 USA 202/653-1034

Andrew R. Dickens

9812 Summerday Drive Burke, VA 22015 USA 203/323-6299

William A. Diener

Jet Propulsion Laboratory 4800 Oak Grove Drive Pasadona, CA 91109 USA 818/354-6670

Rob Douglas

National Research Council of Canada M-36 K1A OR6 Ottawa CANADA 613-993-5186

James M. Eler

U.S. Naval Observatory 3450 Massachusetts Avenue, Northwest Washington, DC 20392-5420 USA 202/653-0350

Donald A. Emmons

Frequency and Time Systems 34 Tozer Road Beverly, MA 01519 USA 508/927-8220

John W. Evans

DISA 10701 Parkridge Boulevard Reston, VA 22091 USA 703/487-3173

Sheila Faulkner

U. S. Naval Observatory 3450 Massachusetts Avenue, Northwest Washington, DC 20392-5420 USA 202/653-1460

Raymond L. Filler

U.S. Army Research AMSRL EP ME Fort Monmouth, NJ 07703-5601 USA 908/544-2467

Richard H. Flamm

Lockheed Missiles and Space Company 1992 Adams Avenue Melbourne, FL 32935 USA 407/867-3145

Henry F. Fliegel

Aerospace Corporation 2350 El Segundo Boulevard El Segundo, CA USA 310/336-1710

Earl Fossler

TRAK Microwave Corporation Systems Division 4726 Eisenhower Boulevard Tampa, FL 33634 USA 813/884-1411

Hugo Fruehauf

Ball Corporation 3 Parker Irvine, CA 92718 USA 714/770-5000

Michael J. Full

NAPA-GPS 955 L'Enfant Plaza North, Southwest Washington, DC 20024 USA 202/651-8014

Ivan J. Galysh

U.S. Naval Research Laboratory Code 8150 4555 Overlook Avenue, Southwest Washington, DC 20375-5000 USA 202/404-7060

Joseph M. Gardner

Leitch, Incorporated 920 Corporate Lane Chesapeake, VA 23302 USA

Michael Garvey

Frequency and Time Systems 34 Tozer Road Beverly, MA 01915 USA 508/927-8220

Richard Gast

TRAK Microwave Corporation Systems Division 4726 Eisenhower Boulevard Tampa, FL 33634 USA 813/884-1411

Glen G. Gibbons

GPS World 859 Willamette Street Eugene, OR 97401 USA 503/984-5235

Al Gifford

U.S. Naval Research Laboratory 1101 Maryland Avenue, Northeast Washington, DC 20002 USA 202/404-5039

Asbjorn M. Gjelsvik

MITRE Burlington Road Bedford, MA 01730 USA 617/271-3712

Raymond L. Granata

NASA/Goddard Space Flight Center Greenbelt, Maryland 20771 USA 301/286-3702

Joe C. M. Green

AlliedSignal Technical Services Corporation 129 North Hill Avenue Pasadena, CA USA 818/584-4472

Charles A. Greenhall

Jet Propulsion Laboratory 4800 Oak Grove Drive MS 298-100 Pasadena, CA 91109 USA 818/393-6944

Michael Grimm

Jet Propulsion Laboratory 4800 Oak Grove Drive Pasadena, CA 91109 USA 818/354-2245

Fran Groat

Hewlett-Packard Company 5301 Stevens Creek Boulevard Santa Clara, CA 95052-8059 USA 408/553-2307

Donald M. Haddox

U.S. Naval Observatory 3450 Massachusetts Avenue, Northwest Washington, DC 20392-5420 USA

Darrell R. Hale

AlliedSignal Technical Services Corporation One Bendix Road MS C-1300 Columbia, MD 21045 USA 410/964-7097

Robert Hamell

Jet Propulsion Laboratory 4800 Oak Grove Drive Pasadena, CA 91109 USA 818/354-4944

Wayne Hanson

National Institute of Standards and Technology 325 Broadway Boulder, CO 80303 USA 303/497-5233

Michael D. Harkins

Applied Technology Associates, Incorporated 6800 Backlick Road Suite 303 Springfield, VA 22150 USA 703/644-0301

Bill Harron

Datum, Incorporated 4545 Pilgrim Lane Boothwyn, PA 19061 USA 610/485-9506

Gregory E. Hatten

U.S. Air Force 2nd Space Operations Squadron 300 O'Malley Avenue Suite 41 Falcon Air Force Base, CO 80912-3041 USA 719/550-6377

Helmut Hellwig

Air Force Office of Scientific Research 110 Duncan Avenue, Suite B115 Bolling Air Force Base Washington, DC 20332-0001 USA 202/767-5017

Thomas P. Holden

Stanford Telecommunications, Incorporated 1221 Crossman Avenue Sunnyvale, CA 94089-1117 USA 408/745-0818

Terry A. Howell

U.S. Naval Observatory 3450 Massachusetts Avenue, Northwest Washington, DC 20392-5420 USA 202/653-1449

Steven T. Hutsell

U.S. Air Force 2nd Space Operations Squadron 300 O'Malley Avenue Falcon Air Force Base, CO 80912-3041 USA 719/550-6394

Jeffrey S. Ingold

AlliedSignal Technical Services Corporation One Bendix Road Columbia, MD 21045 USA 410/964-7188

Andrew Johnson

U.S. Naval Observatory 3450 Massachusetts Avenue, Northwest Washington, DC 20392-5420 USA 202/653-1561

Jeffrey A. Johnson

Overlook Systems Technologies, Incorporated 1950 Old Gallows Road Suite 400 Vienna, VA 22182 USA 703/893-1411

Edward C. Jones

U.S. Naval Research Laboratory Code 8150 4555 Overlook Avenue, Southwest Washington, DC 20375-5000 USA 202/767-0590

Sarunas K. Karuza

The Aerospace Corporation P.O. Box 92957 Los Angeles, CA 90009-2957 USA 213/336-6837

Shalom Kattan

Guide Technology, Incorporated 920 Saratoga Avenue, Suite 205 San Jose, CA 95129 USA 408/246-9903

Robert H. Kern

Kernco, Incorporated 28 Harbor Street Danvers, MA 01923 USA 508/777-1956

Wendy L. King

U.S. Naval Observatory 3450 Massachusetts Avenue, Northwest Washington, DC 20392-5420 USA 202/653-0486

William J. Klepczynski

U.S. Naval Observatory 3450 Massachusetts Avenue, Northwest Washington, DC 20392-5420 USA 202/653-1521

Douglas E. Koch

Naval Space Command 808 Pleasant Hill Lane Bowie, MD 20716 USA 301/249-7494

Paul A. Koppang

U.S. Naval Observatory 3450 Massachusetts Avenue, Northwest Washington, DC 20392-5420 USA 202/653-0350

Guenter Kramer

Physikalish-Technische Bundesanstalt P.O. Box 3345 D-38023 Braunschweig, GERMANY 49 431 592 4420

Paul F. Kuhnle

Jet Propulsion Laboratory 4800 Oak Grove Drive Pasadena, CA 91109 USA 818/354-2715

Eugen F. Kunzi

SIEMENS AG.
Department SI E SY 4
Landshuter Strabe 26
D-85716 Unterschleissheim
GERMANY
49 89 3179 2297

Jack Kusters

Hewlett-Packard Company 5301 Stevens Creek Boulevard Santa Clara, CA 95052-8059 USA 408/553-2041

Graybill P. Landis

U.S. Naval Research Laboratory Code 8150 4555 Overlook Avenue, Southwest Washington, DC 20375-5000 USA 202/404-7061

Marie M. Largay

U.S. Naval Research Laboratory Code 8153 4555 Overlook Avenue, Southwest Washington, DC 20375-5000 USA 202/767-9133

Julius Law

Jet Propulsion Laboratory 4800 Oak Grove Drive Pasadena, CA 91109 USA 818/354-2988

Ming C. Lee

Western Range U.S. Air Force 826 13th Street Vandenberg Air Force Base, CA 93437 USA

Albert Leong

The Aerospace Corporation P.O. Box 92957 MS M1/135 Los Angeles, CA 90009 USA 310-336-6444

Sigfrido M. Leschiutta

Politecnico di Torino 24 Corso Abruzzi 10125 Torino ITALY 38 11 381 8715

Judah Levine

National Institute of Standards and Technology Code 847 325 Broadway Boulder, CO 80303 USA 303/497-3903

Włodzimierz W. Lewandowski

Bureau International des Poids et Mesures Pavillon de Breteuil 92310 Sevres FRANCE 33 1 45 077063

Chuck Little

Hewlett Packard Company 5301 Stevens Creek Boulevard Santa Clara, CA 95052-8059 USA 408/553-2506

Michael A. Lombardi

National Institute of Standards and Technology 325 Broadway Division 847 Boulder, CO 80303 USA 303/497-3212

Gene E. Long

Odetics P.O. Box 2727 LaPlata, MD 20646 USA 301/870-3311

Pete R. Lopez

TRAK Microwave Corporation Systems Division 4726 Eisenhower Boulevard Tampa, FL 33634 USA 813/884-1411

John McK. Luck

Orroral Observatory, Auslis P.O. Box 2 2611 Belconnen ACT AUSTRALIA 61 6 288-3263

Carl F. Lukac

U.S. Naval Observatory 3450 Massachusetts Avenue, Northwest Washington, DC 20392-5420 USA 202/653-1527

Edward M. Lukacs

U.S. Naval Observatory 11820 Southwest 166th Street Miami, FL 33177-2175 USA 305/235-0515

Collin C. MacMillan

Sigma Tau Corporation 1711 Holt Road Tuscaloosa, AL 35404 USA 205/553-0038

Phu V. Mai

U.S. Naval Observatory 3450 Massachusetts Avenue, Northwest Washington, DC 20392-5420 USA 202/653-0350

Lute Maleki

Jet Propulsion Laboratory 4800 Oak Grove Drive Pasadena, CA 91109 USA 818/354-3688

David G. Markham

U.S. Navy COMSPAWARSYSCOM/PMW175 2451 Crystal Drive Arlington, VA 22245 USA

Demetrios N. Matsakis

U.S. Naval Observatory 3450 Massachusetts Avenue, Northwest Washington, DC 20392-5420 USA 202/653-0585

Thomas B. McCaskill

U.S. Naval Research Laboratory Code 8153 4555 Overlook Avenue, Southwest Washington, DC 20375-5000 USA 202/404-7068

Angela Davis McKinley

U.S. Naval Observatory 3450 Massachusetts Avenue, Northwest Washington, DC 20392-5420 USA 202/653-1528

Norma G. Meyer

U.S. Naval Observatory 3450 Massachusetts Avenue, Northwest Washington, DC 20392-5420 USA 202/653-1525

Mihran Miranian

U.S. Naval Observatory 3450 Massachusetts Avenue, Northwest Washington, DC 20392-5420 USA 202/653-1522

Donald H. Mitchell

TrueTime, Incorporated 2835 Duke Court Santa Rosa, CA 95407 USA 707/528-1230

Raymond W. Moskaluk

Hewlett-Packard Company 5301 Stevens Creek Boulevard Santa Clara, CA 95128 USA 408/553-3319

David G. Munton

University of Texas Applied Research Laboratory 10000 Burnet Road Austin, TX 78758 USA 512/835-3831

William J. Murphy

Computer Sciences Corporation P.O. Box 446 Building 1440 Edwards Air Force Base, CA 93523 USA 805/277-2004

James A. Murray

SEA, Incorporated 6215 Houston Street Alexandria, VA 22310 USA 202/404-7057

Craig W. Nelson

SpectraDynamics, Incorporated 445 South 43rd Street Boulder, CO 80303 USA

Clyde C. Norris

Computer Sciences Corporation P. O. Box 217 Clearfield, UT 84015 USA 801/777-4845

Jerry R. Norton

Johns Hopkins University Applied Physics Laboratory Johns Hopkins Road Laurel, MD 21104 USA

Timothy W. Oakley

Sigma Tau Corporation 1711 Holt Road Tuscaloosa, AL 35404 USA 205/553-0038

Jay Oaks

U.S. Naval Research Laboratory Code 8150 4555 Overlook Avenue, Southwest Washington, DC 20375-5000 USA 202/767-1434

David B. Opie

Physical Sciences, Incorporated 635 Slaters Lane Suite G101 Alexandria, VA 22314 USA 703/548-6410

Jorge Osorio

Observatorio Naval Buenos Aires 2099 Av. Espana 1107 Buenos Aires ARGENTINA 64 1 36 1162

Terry N. Osterdock

Stellar GPS Corporation 800 Charcot Avenue Suite 110 San Jose, CA 95131 USA 408/383-1520

H. Bryan Owings

Sigma Tau Corporation 1711 Holt Road Tuscaloosa, AL 35404 USA 205/553-0038

Thomas E. Parker

National Institute of Standards and Technology 325 Broadway Boulder, CO 80303 USA 303/497-7881

Ralph E. Partridge

Los Alamos National Laboratory DX-12, MS-P947 P.O. Box 1663 Los Alamos, NM 87545 USA 505/667-5255

William H. Paul

Tecom NRO-DT-T White Sands Missile Range, NM 88002 USA 505/674-9220

Peter Z. Paulovich

NISE East Det Norfolk P. O. Box 1376 Norfolk, VA 23501-1376 USA 804/396-0287

Bruce Penrod

TrueTime, Incorporated 2835 Duke Court Santa Rosa, CA 95407 USA 707/528-1230

James C. Perry

NASA/Goddard Space Flight Center Greenbelt, MD 20771 USA 301/286-3471

Harry E. Peters

Sigma Tau Standards Corporation 1711 Holt Road Tuscaloosa, AL 35404 USA 205/553-0038

Larry E. Peters

Sigma Tau Standards Corporation P.O. Box 1877 Tuscaloosa, AL 35403 USA 205/553-0038

James S. Phelps

Scientific and Technical Analysis 11300 Naples Mill Road Fairfax, VA 22043 USA 703/934-0188

William D. Phillips

National Institute of Standards and Technology PHY A167 Gaithersburg, MD 20899 USA 301/975-6554

Eva F. Pikal

National Institute of Standards and Technology 325 Broadway Boulder, CO 80303 USA 303/497-5629

William M. Powell

U. S. Naval Observatory 3450 Massachusetts Avenue, Northwest Washington, DC 20392-5420 USA 202/653-1528

Edward D. Powers

U.S. Naval Research Laboratory Code 8150 4555 Overlook Avenue, Southwest Washington, DC 20375-5000 USA 202/767-5004

Robert E. Price

AlliedSignal Technical Services Corporation One Bendix Road Columbia, MD 21045 USA 410/964-7437

David W. Rea

Spectrum Geophysical Instruments 1900 West Garvey Avenue, South Suite 200 West Covina, CA 91790 USA 714/544-3000

Elza K. Redman

DoD 9800 Savage Road Fort Meade, MD 20577 USA 301/688-7526

Wilson G. Reid

Naval Research Laboratory Code 8150 4555 Overlook Avenue, Southwest Washington, DC 20375-5000 USA 202/404-7068

Victor S. Reinhardt

Hughes Space and Communications P.O. Box 92919 Los Angeles, CA 90009 USA

William J. Riley

EG&G Rubidium 35 Congress Street Salem, MA 01970 USA 508/745-3200

Eric B. Rodal

Trimble Navigation 585 North Mary Sunnyvale, CA 94088 USA 408/481-2039

Ronald C. Roloff

Datum, FTS 8005 McKenstry Drive Laurel, MD 20723 USA 301/725-3636

Harry W. Sadler III

AlliedSignal Technical Services Corporation One Bendix Road Columbia, MD 21045 USA 410/964-7425

Roberto O. Sartori

Observatorio Naval Buenos Aires 2099 Av. Espana 1107 Buenos Aircs ARGENTINA 54 1 361 1162

Bernard R. Schlueter

Observatory of Neuchatel CH2000 Neuchatel SWITZERLAND 41 38 241861

Richard E. Schmidt

U.S. Naval Observatory 3450 Massachusetts Avenue, Northwest Washington, DC 20392-5420 USA 202/653-0487

Steve Schumacher

Loral Space and Range Systems 2900 Murrell Road Rockledge, FL 32955 USA 407/635-7647

Mike Shunfenthal

TASC 12100 Sunset Hills Road Reston, VA 22090 USA 703/834-5000

Ralph E. Simons

Frequency and Time Systems 34 Tozer Road Beverly, MA 01915 USA 508/927-8220

Hank T. Skalski

DOT/OST 400 7th Street, Southwest Washington, DC 20590 USA 202/366-4894

Luke R. Smith

Sigma Tau Standards Corporation 1711 Holt Road Tuscaloosa, AL 35403 USA 205/553-0038

Benjamin F. Smith

Stanford Telecommunications, Incorporated 1221 Crossman Avenue Sunnyvale, CA 94089-1117 USA 408/745-2640

Andrew G. Snow

Frequency and Time Systems 34 Tozer Road Beverly, MA 01915 USA 508/927-8220

Armin Soering

Deutsche Telekom AG 3 AM Kavalleriesand 64295 Darmstadt GERMANY 49 6151 83 4549

Samuel R. Stein

Timing Solutions Corporation 1025 Rosewood Avenue, #200 Boulder, CO 80304 USA 303/939-8481

Michael A. Street

Bonneville Power Administration 905 Northeast 11th Avenue Portland, OR 92232 USA 503/230-4363

Richard L. Sydnor

Jet Propulsion Laboratory 4800 Oak Grove Drive Pasadena, CA 91109 USA 818/354-2763

Philip E. Talley, Jr.

Retired 1022 Eagle Crest Macon, GA 31211 USA 913/745-4661

Claudine Thomas

Bureau International des Poids et Mesures Pavillon de Breteuil 92312 Sevres Cedex FRANCE 33 1 45 07 7073

Larry M. Thomas

Computer Sciences Corporation 44609 North 86th Street East Lancaster, CA 93535 USA 805/277-6130

Roger Thornburn

Stellar GPS Corporation 800 Charcot Avenue Suite 110 San Jose, CA 95131 USA 408/383-1522

Terry L. Trisler

Ball Corporation Efratom Division 3 Parker Irvine, CA 92718 USA 714/770-5000

Gregory B. Turetzky

Trimble Navigation 645 North Mary Avenue Sunnyvale, CA 94086 USA 408/252-9510

David A. Turner

National Research Council 2101 Constitution Avenue, Northwest HA-286 Washington, DC USA 202/334-2895

Pierre J. M. Uhrich

LPTF Observatoire de Paris 61 Avenue de L'Observatoire F-75014 Paris FRANCE 33 1 40 512216

M. J. VanMelle

Rockwell 5505 Flintridge Drive Colorado Springs, CO 80918 USA 719/550-2705

Francine M. Vannicola

U.S. Naval Observatory 3450 Massachusetts Avenue, Northwest Washington, DC 20392-5420 USA 202/653-1525

Vince Vannicola

U.S. Coast Guard 7323 Telegraph Road Alexandria, VA 22015 USA 703/313-5811

Christian Veillet

Observatoire de La Côte d'Azur Avenue Nicolas Copernic 06130 Grasse FRANCE

John R. Vig

U.S. Army Research Laboratory AMSRL-EP-ME Fort Monmouth, NJ 07703-5601 USA

Frank J. Voit

The Aerospace Corporation P.O. Box 92957 Los Angeles, CA 90009-2957 USA 310/336-6764

Warren F. Walls

Femtosecond Systems P.O. Box 6005 Columbia, MD 21045-8005 USA 410/740-1427

Heidi E. Walls

Femtosecond Systems P.O. Box 6005 Columbia, MD 21045-8005 USA 410/740-1427

Fred L. Walls

National Institute of Standards and Technology 325 Broadway Boulder, CO 80303 USA 303/497-3207

Harry T. M. Wang

Hughes Research Lab 3011 Malibu Canyon Road Malibu, CA 90265 USA 805/317-5431

Fred L. Walls

National Institute of Standards and Technology 325 Broadway Boulder, CO 80303 USA 303/497-3207

Harry T. M. Wang

Hughes Research Lab 3011 Malibu Canyon Road Malibu, CA 90265 USA 805/317-5431

S. Clark Wardrip

AlliedSignal Technical Services Corporation 726 Foxenwood Drive Santa Maria, CA 93455-4221 USA 805/937-6448

Marc A. Weiss

National Institute of Standards and Technology 325 Broadway Boulder, CO 80303 USA 303/497-3261

Paul J. Wheeler

U.S. Naval Observatory 3450 Massachusetts Avenue, Northwest Washington, DC 20392-5420 USA 202/653-0516

Joseph D. White

U.S. Naval Research Laboratory Code 8151 4555 Overlook Avenue, Southwest Washington, DC 20375-5000 USA 202/767-5111

Carol A. Williams

Department of Mathematics University of South Florida 4202 East Fowler Avenue PHY 114 Tampa, FL 33620 USA 813/974-3553

Warren L. Wilson

Lockheed Missiles and Space Company 701 East Thrift Avenue Kingsland, GA 31548-8213 USA 912/673-1148

Gernot M. R. Winkler

U. S. Naval Observatory 3450 Massachusetts Avenue, Northwest Washington, DC 20392-5420 USA 202/653-1520

Peter Wolf

Bureau International des Poids et Mesures Pavillon de Breteuil 92312 Sevres Cedex FRANCE 33 14 507 7075

William H. Wooden

Defense Mapping Agency 8613 Lee Highway Stop A-13 Fairfax, VA 22031-2137 USA 703/285-9339

David F. Wright

Radiocode Clocks Ltd. 54 Watling Street South 5467BQ Church Stretton Shropshire UNITED KINGDOM 44 1694 723691

James L. Wright

Computer Sciences Raytheon P.O. Box 4127 Patrick Air Force Base, FL 32925 USA 407/494-2014

Andy Wu

The Aerospace Corporation 2350 East El Segundo El Segundo, CA 90245-4691 USA 310/336-0437

Victor S. Zhang

National Institute of Standards and Technology 325 Broadway Boulder, CO 80303 USA 303/497-3977

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REPORT DOCUMENTATION PAGE

Form Approved OMB No. 0704-0188

Public reporting burden for this collection of information is estimated to average 1 hour per response, including the time for reviewing instructions, searching existing data sources, gathering and maintaining the data needed, and completing and reviewing the collection of information. Send comments regarding this burden estimate or any other aspect of this collection of information, including suggestions for reducing this burden, to Washington Headquarters Services, Directorate for Information Operations and Reports, 1215 Jefferson Davis Highway, Suite

1. AGENCY USE ONLY (Leave blank)	2. REPORT DATE May 1995	3. REPORT TYPE AND DATES COVERED Conference Publication	
4. TITLE AND SUBTITLE 26th Annual Precise Time and Time Interval (PTTI) Applications and Planning Meeting		5. FUNDING NUMB Code 50	
6. AUTHOR(S) Dr. Richard Sydnor, Techn	nical Editor		
7. PERFORMING ORGANIZATION NAMI Goddard Space Flight Cen Greenbelt, Maryland 2077	ter	8. PERFORMING O	
		95B000	83
9. SPONSORING/MONITORING AGENCY NAME(S) AND ADDRESS(ES) NASA Aeronautics and Space Administration		10. SPONSORING/ AGENCY REP	
Washington, D.C. 20546-0001		CP-330	2
1. SUPPLEMENTARY NOTES		l	
Laboratory; Space and Naval Warfa Devices Laboratory; Rome Laboratory	re Systems Command; Na orv: and Air Force Office		Fechnology and
Unclassified-Unlimited Subject Category: 70 Report available from the NASA Landing Road, Linthicum Height	Center for AeroSpace Inf		N CODE

This document is a compilation of technical papers presented at the 26th Annual PTTI Applications and Planning Meeting, held December 6 through December 8, 1994, at the Hyatt Regency Hotel, Reston, Virginia. Papers are in the following categories:

- Recent developments in rubidium, cesium, and hydrogen-based frequency standards, and in cryogenic and trappedion technology.
- International and transnational applications of Precise Time and Time Interval technology with emphasis on satellite laser tracking, GLONASS timing, intercomparison of national time scales and international telecommuni-
- Applications of Precise Time and Time Interval technology to the telecommunications, power distribution, platform positioning, and geophysical survey industries.
- Applications of PTTI technology to evolving military communications and navigation systems.
- Dissemination of precise time and frequency by means of GPS, GLONASS, MILSTAR, LORAN, and synchronous communications satellites.

14. SUBJECT TERMS			15. NUMBER OF PAGES
Frequency Standards, Hydroge Synchronization, Precise Time Phase Noise	505 16. PRICE CODE		
17. SECURITY CLASSIFICATION OF REPORT	18. SECURITY CLASSIFICATION OF THIS PAGE	19. SECURITY CLASSIFICATION OF ABSTRACT	20. LIMITATION OF ABSTRACT
Unclassified	Unclassified	Unclassified	Unlimited